

# Enhanced Electromechanical Modeling of Asymmetrical Dual Three-Phase IPMSM Drives

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**Abstract**—Throughout literature, a clear concern for achieving high performance stator current control of asymmetrical dual three-phase IPMSM drives as well as fault tolerant strategies were demonstrated. However, limited attention was given on how to effectively compute, simulate and compare the expected (i.e. simulated) and the practically applied electromechanical torque. This requires an accurate electromechanical model with precise electromechanical torque estimation. This paper presents a high fidelity *direct-quadrature* ( $dq$ ) model capable of describing the generated electromechanical torque. The presented method is based on obtaining the non-linear flux linkage maps as well as describing the harmonic models, which contribute to the average and ripple torque, respectively. The flux linkage maps contribute in the tuning of the stator current controllers and electromechanical torque computation, which is compared to the computed torque from the manufacturer parameters and literature. The proposed model does not require previous knowledge concerning geometric and machine design data. The results coincide to a high extent with the experimental validations on a 2.5 kW asymmetrical dual three-phase IPMSM prototype.

**Index Terms**—Current control, Dual three-phase, IPMSM, modeling.

## NOMENCLATURE

$\mathbb{R}, \mathbb{N}$  are the set of real and natural numbers.  $x \in \mathbb{R}$  is a real scalar, while  $\mathbf{x} \in \mathbb{R}^n$  (bold) is a real valued vector with  $n \in \mathbb{N}$ .  $\mathbf{x}^T$  is the transpose of  $\mathbf{x}$ .  $\mathbf{X} \in \mathbb{R}^{n \times m}$  (capital bold) is a real valued matrix with  $n \in \mathbb{N}$  rows and  $m \in \mathbb{N}$  columns.  $\mathbf{I}_n \in \mathbb{R}^{n \times n} := \text{diag}(1, \dots, 1)$ : identity matrix.  $\mathbf{O}_n \in \mathbb{R}^{n \times n} := \text{diag}(0, \dots, 0)$ : zero square matrix. The determinant of a square matrix  $\mathbf{Q}$  is  $\det(\mathbf{Q})$ .  $\mathbf{x}_s^{a_1 \rightarrow c_2} \in \mathbb{R}^6$  is the stator space vector expressed in the  $(a_1 b_1 c_1 - a_2 b_2 c_2)$  reference frame and may represent voltage  $\mathbf{u}$  (V), flux linkage  $\boldsymbol{\psi}$  (Wb) or current  $\mathbf{i}$  (A), i.e.  $\mathbf{x} \in \{\mathbf{u}, \boldsymbol{\psi}, \mathbf{i}\}$ . The subscript 's' denotes stator quantities. The electromechanical and load torques ( $\text{N} \cdot \text{m}$ ), the electrical angular speed (rad/s), electrical angular position (rad) with respect to the reference flux axis of phase  $a_1$ , and inertia ( $\text{kg} \cdot \text{m}^2$ ) are  $m_e$ ,  $m_{\text{load}}$ ,  $\omega_e$ ,  $\phi_e$ , and  $\Theta$ , respectively. The viscous friction coefficient is  $\nu$  ( $\text{N} \cdot \text{m} \cdot \text{s}$ ), while the coulomb friction is  $m_c$  ( $\text{N} \cdot \text{m}$ ). The electrical resistance ( $\Omega$ ) and inductances (H) are indicated by  $R$  and  $L$ , respectively. Sampling / switching frequencies is given by  $f_{\text{sw}}$  (kHz).

## I. INTRODUCTION

The rapid development of power electronic devices and real-time embedded systems promoted significantly the adoption of multiphase (i.e. more than three-phases) drives, which are employed when a high level of fault-tolerance and reliability are of interest [1]. Within the last decade, a noticeable amount of literature has been dedicated to dual three-phase (DTM) drives; since dual three-phase machines can be easily realized through rewinding of the existing three-phase stator windings [2]. Such rewinding permits varying the space-orientation of the two three-phase sets with respect to each other. Most popular space angles are *zero*, also known as dual-windings,  $\pi/3$  rad symmetrical dual three-phase windings, and  $\pi/6$  rad known as asymmetrical dual three phase (ADT) windings [3]. The asymmetrical type has better fundamental flux content over the symmetrical type, but slightly lower fault-tolerance capability [3], [4]. However, this paper deals with the asymmetrical type. Asymmetrical dual three-phase IPMSMs (ADT-IPMSMs) are known for their high efficiency and improved power/weight ratio, when compared to ADT induction machines (ADT-IMs), making ADT-IPMSMs potential candidates for instance for currently emerging technologies as electric vehicles [5].

Zhao et. al. in [6] initiated the principle of harmonics distribution in the different control subspaces for ADT-IMs, where such harmonics are mainly due to the converter and windings non-idealities. For ADT-IPMSMs, the developed harmonics are mainly due to the permanent magnet (PM) non-sinusoidal distribution in the air-gap. Compensating such harmonics is essential towards achieving high-performance operation. Lately, research developed and excelled in achieving a high performance operation of ADT-IPMSMs through realizing, to a high extent, a sinusoidal harmonic-free stator currents. This was accomplished by means of anti-, synchronous proportional-integral-resonant (PIR) controllers [7], multi-resonant controllers [8], and disturbance observers [9]. However, to the best of the authors' knowledge, not much attention has been reported on how to obtain a high fidelity electromechanical model for ADT-IPMSMs in the torque-producing  $dq$  frame, which takes into account the fact that IPMSMs are practically *non-linear* machines. Such non-linear model is capable of describing effectively the rated electromechanical torque as well as anticipating the electrical transients.

This paper demonstrates a mathematical model capable of generating both: the fundamental and harmonic elements in the  $dq$  frame of ADT-IPMSMs, which are practically *non-linear* machines. The generated electromechanical torque is a function of both the fundamental and harmonic elements. The fundamental dynamics are identified by means of the experimentally identified fundamental flux linkage maps, while the harmonics are generated by means of the experimentally identified harmonic flux linkage parameters through an open circuit test. The cross-coupling between the fundamental  $dq$  subspace flux-linkages (i.e. due to saturation) is taken into account. The fundamental flux linkage maps are adopted in designing the  $dq$  current controller. The entire experiments and conclusions rely on the stator current measurements and encoder feedback, which is a standard in drive systems, without prior knowledge of the geometric machine design or finite element data. It is also concluded that taking into account the ADT-IPMSM non-linearity increases the expected torque density of the machine and enhances the operation on the Maximum-Torque-per-Ampere (MTPA) current loci [10]–[12]. Using the flux linkage maps, the computed electromechanical torque is compared to the computed torque from the manufacturer parameters and literature, at the same conditions, showing the superiority of the proposed method. A 2.5 kW ADT-IPMSM experimental setup was employed for validation.

## II. ASYMMETRICAL DUAL THREE PHASE DRIVE MODEL

### A. Dual three phase IPMSM model

For a DTM double-star neutral drive, a generalized Clarke's transformation, also known as vector space decomposition (VSD) [3], is given by

$$\mathbf{T}_{\text{VSD}} = \frac{1}{3} \begin{bmatrix} 1 & -\frac{1}{2} & -\frac{1}{2} & \frac{\sqrt{3}}{2} & -\frac{\sqrt{3}}{2} & 0 \\ 0 & \frac{\sqrt{3}}{2} & -\frac{\sqrt{3}}{2} & \frac{1}{2} & \frac{1}{2} & -1 \\ 1 & -\frac{1}{2} & -\frac{1}{2} & -\frac{\sqrt{3}}{2} & \frac{\sqrt{3}}{2} & 0 \\ 0 & -\frac{\sqrt{3}}{2} & \frac{\sqrt{3}}{2} & \frac{1}{2} & \frac{1}{2} & -1 \end{bmatrix}, \quad (1)$$

which decomposes the asymmetrical six-phases quantities  $\mathbf{f}_s^{a_1 \rightarrow c_2}$ , where  $f \in \{u, \psi, i\}$ , to two orthogonal planes  $(\mathbf{f}_s^{\alpha\beta}, \mathbf{f}_s^{xy}) = \mathbf{T}_{\text{VSD}} \mathbf{f}_s^{a_1 \rightarrow c_2}$ , namely the equivalent  $\alpha\beta$  subspace and the  $xy$  subspace, which controls the degree of unbalance between the three-phase sets and derating factors during faults [1], [3], [13]. The VSD matrix maps different harmonics to such subspaces, where  $\alpha\beta$  plane holds the  $12\gamma \pm 1$  harmonics, where  $\gamma = \{1, 2, 3, \dots\}$ ,  $xy$  is concerned with the  $6\Gamma \pm 1$  harmonics, where  $\Gamma = \{1, 3, 5, \dots\}$ . Using the generalized Park's transformation [7]

$$\bar{\mathbf{T}}_p(\phi_e)^{-1} = \begin{bmatrix} \mathbf{T}_p(\phi_e)^{-1} & \mathbf{O}_2 \\ \mathbf{O}_2 & \mathbf{I}_2 \end{bmatrix}, \quad (2)$$

where  $\mathbf{T}_p(\phi_e)^{-1} = \begin{bmatrix} \cos(\phi_e) & -\sin(\phi_e) \\ \sin(\phi_e) & \cos(\phi_e) \end{bmatrix}$ , the ADTM-IPMSM dynamic model is

TABLE I. Parameters of the ADT-IPMSM drive.

Parameter	Value
Stator resistance	$R_s = 1.1 \Omega$
Nominal d-axis inductance*	2 mH
Nominal q-axis inductance*	11.5 mH
Computed d-axis inductance**	6.7 mH
Computed q-axis inductance**	20 mH
Leakage inductance	$L_l = 0.9 \text{ mH}$
Pole-pair	$n_p = 3$
Inertia	$\Theta = 0.01 \text{ kg}\cdot\text{m}^2$
Viscous friction	$\nu = 0.8 \times 10^{-3} \text{ N}\cdot\text{m}\cdot\text{s}$
Coulomb friction	$m_c = 0.2725 \text{ N}\cdot\text{m}$
Rated stator current	$i_{s,\text{rated}} = 4.2 \text{ A}$
Rated torque	$m_{e,\text{rated}} = 10.6 \text{ N}\cdot\text{m}$
Rated mechanical speed	2300 RPM
DC-link voltage	$u_{dc} = 580 \text{ V}$
Sampling and switching rates	$f_{sw} = 8 \text{ kHz}$

\* Half the manufacturer's value of the original three-phase machine due to halving the number of turns when carrying-out the rewinding to the dual three-phase machine.

\*\* Computed as shown in [15, Sec. 5.3].

$$\left. \begin{aligned} \bar{\mathbf{u}}_s^{dq} &= R_s \bar{\mathbf{i}}_s^{dq} + \omega_e \mathbf{J} \bar{\boldsymbol{\psi}}_s^{dq} + \frac{d}{dt} \bar{\boldsymbol{\psi}}_s^{dq}, \\ \bar{\mathbf{u}}_s^{xy} &= R_s \bar{\mathbf{i}}_s^{xy} + L_l \frac{d}{dt} \bar{\mathbf{i}}_s^{xy}, \\ \frac{d}{dt} \omega_e &= \frac{n_p}{\Theta} (m_e - m_{\text{load}} - \nu \frac{\omega_e}{n_p} - m_c), \end{aligned} \right\} \quad (3)$$

where  $\mathbf{J} := \mathbf{T}_p(\pi/2)^{-1}$ ,  $\bar{\mathbf{i}}_s^{xy} = (i_s^x, i_s^y)^\top$ .  $\bar{\mathbf{f}}_s^{dq} = (\bar{f}_s^d, \bar{f}_s^q)^\top$ , where  $f \in \{\psi, u, i\}$ , is the total (i.e. fundamental and harmonic elements) flux linkage/voltage/current, respectively, in the  $dq$  frame. Note that the  $dq$  and  $xy$  are decoupled subspaces, thus, the  $xy$  plane do not interfere in the torque production mechanism [7], [8], [14]. The ADT-IPMSM parameters are shown in Table I. The model in (3) allows for non-linear flux linkage variation in the  $dq$  frame with the corresponding currents. Thus, in this paper, the  $dq$  flux linkages are used as *states* instead of currents.

### B. $dq$ flux linkage and torque models

The unique difference between different ADTMs is how  $\bar{\boldsymbol{\psi}}_s^{dq}$  in (3) is modelled. Unlike the flux linkage models and the assumed constant  $dq$  inductances presented in [8], [15], [16], the current-dependency of inductance matrices as well as the harmonic flux linkages are considered in the proposed model [20, Chap. 14]. As explained in Sect. II-A, the  $dq$  frame holds the fundamental component, which contributes mainly to the average electromechanical torque and the  $dq$  harmonic fluxes, sourced by the PM harmonics, contribute to the oscillating torque. Accordingly, in order to show the total generated electromechanical torque, it is necessary to have a dynamic model capable of modeling both the fundamental and harmonic components. Thus,  $\bar{\boldsymbol{\psi}}_s^{dq}$  in (3) can be written as

$$\bar{\boldsymbol{\psi}}_s^{dq} = \boldsymbol{\psi}_s^{dq} + \boldsymbol{\psi}_{s,\text{PM}}^{dq}, \quad (4)$$

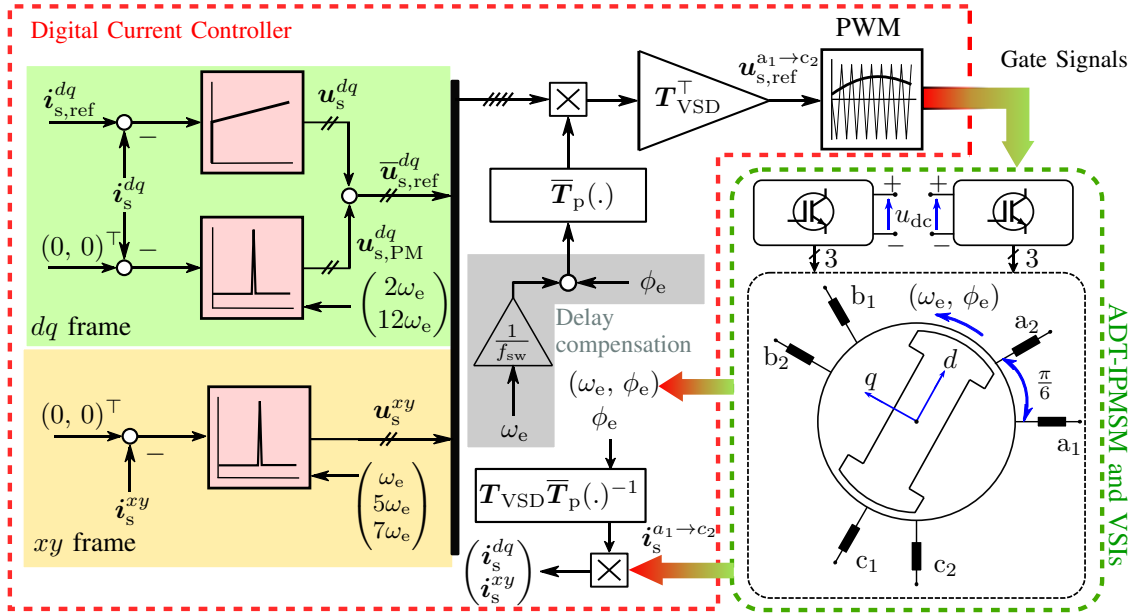


Fig. 1. Proportional-integral and resonant current controllers for the  $dq$  and  $xy$  frames with the corresponding resonant frequencies for an ADT-IPMSM [8].

where,  $\psi_s^{dq}$  is the fundamental flux linkage term, while  $\psi_{s,PM}^{dq}$  is the harmonic term. The same principle is applied to decompose  $\bar{u}_s^{dq}$  into fundamental and harmonic components.

As for  $\bar{i}_s^{dq}$ , the impedance of the  $dq$  frame is assumed to be large enough to neglect  $i_{s,PM}^{dq}$  (i.e.  $\bar{i}_s^{dq} \approx i_s^{dq}$ ). With that approximation in mind, the electromechanical torque produced is denoted by

$$\bar{m}_e = 3n_p i_s^{dq\top} \mathbf{J} (\psi_s^{dq} + \psi_{s,PM}^{dq}) = m_e + m_{e,PM}, \quad (5)$$

where the first term in (5) is the average torque generated by the fundamental component, while the second term represents the oscillatory torque term.

Assuming that non-fundamental components are compensated, as shown later in Sect. III-B, the relation between  $\psi_s^{dq}$  and  $i_s^{dq}$  is *non-linear* with a cross-coupling between the  $dq$  coordinates. This is mathematically modeled as

$$\frac{d}{dt} \psi_s^{dq} = \mathbf{L}_s^{dq} (i_s^{dq}) \frac{d}{dt} i_s^{dq}, \quad (6)$$

where  $\mathbf{L}_s^{dq}(i_s^{dq})$  is the *differential* inductance matrix [20, Chap. 14] and is defined as

$$\mathbf{L}_s^{dq}(i_s^{dq}) := \begin{bmatrix} \frac{\partial \psi_s^d}{\partial i_s^d} & \frac{\partial \psi_s^d}{\partial i_s^q} \\ \frac{\partial \psi_s^q}{\partial i_s^d} & \frac{\partial \psi_s^q}{\partial i_s^q} \end{bmatrix} = \begin{bmatrix} L_s^d & L_m^{dq} \\ L_m^{dq} & L_s^q \end{bmatrix}. \quad (7)$$

This approach is similar to that shown in [17], which is also known as the flux linkage maps and the corresponding differential inductances identification. The current vector  $i_s^{dq}$  is varied to cover the entire operating range of the adopted ADT-IPMSM at a given speed. In the  $xy$  frame, the reference currents are set to  $i_{s,ref}^{xy} = (0, 0)^\top$ , as shown in Fig 1. Finally, for the given  $R_s$  in Table I and a non-zero speed (i.e.  $\omega_e \neq 0$ )

at *steady-state*, the flux linkage vector is obtained by

$$\psi_s^{dq} = \frac{\mathbf{J}^{-1}}{\omega_e} (\mathbf{u}_s^{dq} - R_s i_s^{dq}). \quad (8)$$

The acquired flux linkages and the corresponding differential inductances are shown in Figs. 2(a)-(b) and Figs. 3(a)-(c), respectively.

Since the  $dq$  frame encompasses the rotor 11<sup>th</sup> and 13<sup>th</sup> harmonics, applying Park's transformation in (2) maps such harmonics in the  $dq$  plane, rotating at  $12\omega_e$ . The proposed model for  $\psi_{s,PM}^{dq}$  is

$$\psi_{s,PM}^{dq} = \begin{pmatrix} \psi_{pm}^{13} \sin(12\omega_e t + \phi_o^{13}) - \psi_{pm}^{11} \sin(12\omega_e t + \phi_o^{11}) \\ \psi_{pm}^{13} \cos(12\omega_e t + \phi_o^{13}) + \psi_{pm}^{11} \cos(12\omega_e t + \phi_o^{11}) \end{pmatrix}, \quad (9)$$

where the constants  $\psi_{pm}^{11}$ ,  $\psi_{pm}^{13}$ ,  $\phi_o^{11}$ , and  $\phi_o^{13}$  are identified through the open circuit test as explained in Sect. IV. Finally the equivalent fundamental and harmonic circuit for the  $dq$  frame is shown in Fig. 4.

### III. ADT-IPMSM CURRENT CONTROLLER

#### A. PI current controllers and disturbance compensation

Upon computing the flux linkage maps and the corresponding differential inductances, the PI controllers in Fig. 1 are designed according to the magnitude optimum criterion along with updating the PI controller gain and time constant [17]. Accordingly, the proportional  $V_p$  and integral time constant  $T_i$  of the non-linear PI controllers are defined as

$$\left. \begin{aligned} (V_p^d \ V_p^q)^\top &= \frac{f_{sw}}{2} (\hat{L}_s^d \ \hat{L}_s^q)^\top, \\ (T_i^d \ T_i^q)^\top &= \frac{1}{R_s} (\hat{L}_s^d \ \hat{L}_s^q)^\top, \end{aligned} \right\} \quad (10)$$

where

$$(\hat{L}_s^d \ \hat{L}_s^q)^\top = \det(\mathbf{L}_s^{dq}(i_s^{dq})) \left( \frac{1}{L_s^d} \ \frac{1}{L_s^q} \right)^\top. \quad (11)$$

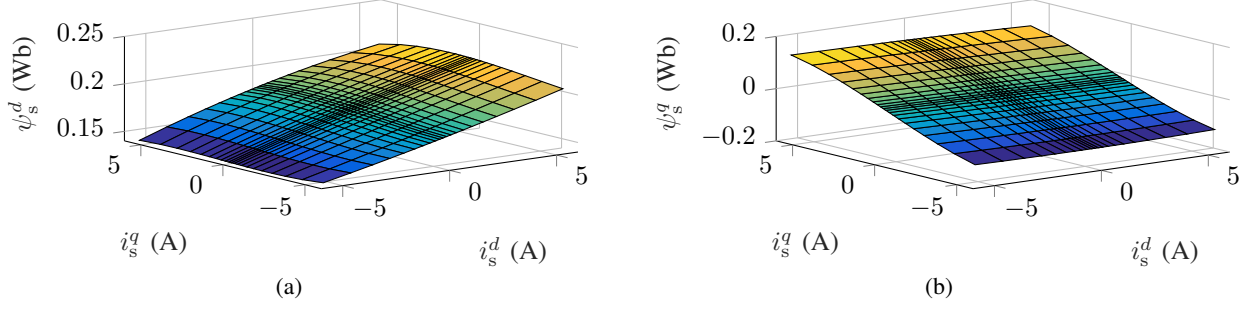


Fig. 2. Measured flux linkage maps of  $\psi_s^{dq}$  in (3), showing the variation of (a)  $\psi_s^d$  and  $\psi_s^q$  for  $\|\mathbf{i}_s^{dq}\| \leq i_{s,rated}$ .

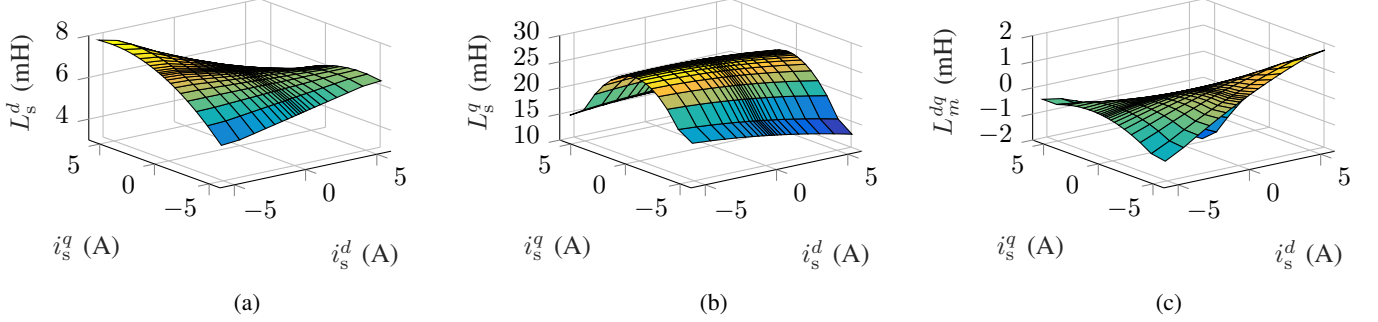


Fig. 3. Measured differential inductances of  $L_s^{dq}(\mathbf{i}_s^{dq})$  for  $\|\mathbf{i}_s^{dq}\| \leq i_{s,rated}$  (left to right): (a)  $L_s^d$ , (b)  $L_s^q$ , and (c)  $L_m^{dq}$ .

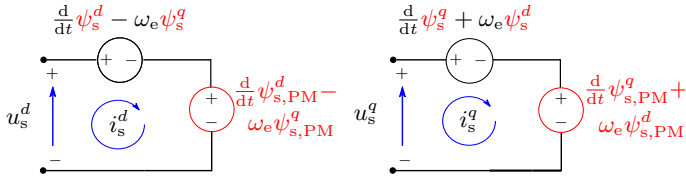


Fig. 4. Proposed equivalent circuits for the  $dq$  frame along with highlighting in red the enhanced model elements as explained in Subsect. II-A.

In [17], the disturbance feedforward compensation enhances the tracking speed of the PI controllers. The disturbance voltage  $\mathbf{u}_{s,dist}^{dq}$  encompasses the coupling terms between the  $dq$  frame coordinates, assuming the machine parameters are identified, which is the case here upon computing  $\psi_s^{dq}$  and  $L_s^{dq}(\mathbf{i}_s^{dq})$ . As proven in [17],  $\mathbf{u}_{s,dist}^{dq}$  is computed by

$$\mathbf{u}_{s,dist}^{dq} = -L_m^{dq} \begin{bmatrix} 0 & \frac{1}{L_s^q} \\ \frac{1}{L_s^d} & 0 \end{bmatrix} (\mathbf{u}_s^{dq} - R_s \mathbf{i}_s^{dq} - \omega_e \mathbf{J} \psi_s^{dq}) - \omega_e \mathbf{J} \psi_s^{dq}. \quad (12)$$

### B. Resonant (R) controllers

As illustrated in Sect. II-A, the presence of harmonic flux linkage elements in the  $dq$  and  $xy$  frames invokes the presence of harmonic currents within such frames. Compensating such current harmonics is inevitable to guarantee high quality performance for the employed ADT-IPMSM drive. This could be achieved either by adopting multi-synchronous frames in both the  $dq$  and  $xy$  frames [7] or by invoking multi-resonant controllers for the shown model in (3) [8]. The

latter method is the one adopted in this paper. Since the amplitude of harmonic currents are speed dependent, meaning they are boosted/curtailed with the increase/decrease of  $\omega_e$ , the resonant (R) controllers are expected to be speed-dependent tuned. The R-controllers are discretized using the impulse-invariant method to ensure precise lock on the desired resonant frequency  $\omega_R$  as well as stable operation [8], [18]. The time constant of such R-controllers  $T_R$  is carried out similar to the method shown in [8] which leads to

$$(T_R^d \ T_R^q \ T_R^{xy}) = \frac{100}{R_s \omega_e} (\hat{L}_s^d \ \hat{L}_s^q \ L_l), \quad (13)$$

where  $(T_R^d, T_R^q)$  are adapted online with  $L_s^d$  and  $L_s^q$ , respectively. The tuned resonant frequencies of the R-controllers in the  $dq$  and  $xy$  frames are  $\omega_R^{dq}$  and  $\omega_R^{xy}$ , respectively. As seen in Fig. 1, there are two pairs of R-controllers in the  $dq$  frame where  $\omega_R^{dq} = (2\omega_e, 12\omega_e)$ , while three pairs of R-controllers in the  $xy$  frame, tuned at  $\omega_R^{xy} = (\omega_e, 5\omega_e, 7\omega_e)$ .

## IV. EVALUATION RESULTS

### A. Experimental environment

This section evaluates the outcomes of the proposed model on a practical setup, shown in Fig. 5, having the same parameters as indicated in Table I. The test bench in Fig. 5 consists of an ADT-IPMSM tied to a standard three-phase induction machine acting as the prime mover. With the aid of Matlab/Simulink, the control algorithm is implemented on dSPACE DS1007. The PWM signals are sent to the VSIs via the DS5101 PWM card, where the triggering signals are manually programmed to ensure precise timings and sampling

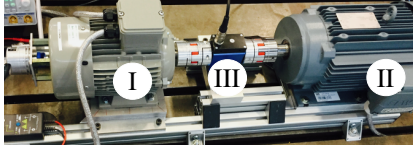


Fig. 5. Test bench: (I) ADT-IPMSM, (II) induction machine (i.e. prime mover) and (III) torque sensor.

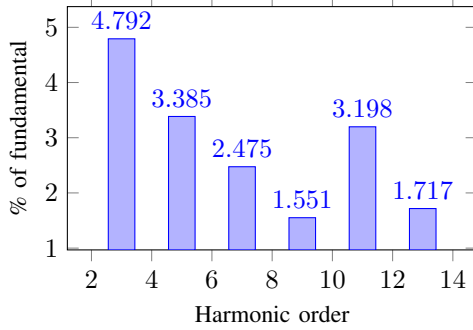


Fig. 6. Back-*e.m.f* harmonic content of the ADT-IPMSM at 1000 RPM.

TABLE II. FFT of the rotor flux linkage at 1000 RPM.

Order	Magnitude (Wb)	Phase (rad)
1	$180 \times 10^{-3}$	0
11	$\psi_{pm}^{11} = 5.1 \times 10^{-3}$	$\phi_o^{11} = 3.5168$
13	$\psi_{pm}^{13} = 2.4 \times 10^{-3}$	$\phi_o^{13} = 0.4224$

instants. The A/D DS2004 board is triggered by the DS5101 board to carry-out the measurements in the middle of the switching period. The real-time control structure is similar to the indicated block diagram in Fig. 1. The discrete-time delay in the position measurement is compensated when applying the Park transformation (i.e.  $\bar{T}_p(\phi_e + \omega_e/f_{sw})$ ), as highlighted in Fig. 1 [19]. To identify the constants in (9), the ADT-IPMSM is open circuited and rotated at a constant speed of 1000 RPM. The back-*e.m.f* was measured using an oscilloscope and analyzed using Fast Fourier Transform (FFT) shown in Fig. 6, from which the flux-linkage harmonic content was estimated as shown in Table II. Moreover, the electromechanical torque is computed on dSPACE three times using the flux linkage maps in (5), obtained parameters in [15, Sec. 5.3], and the manufacturer parameters. The computed electromechanical torques are compared against the practically developed electromechanical torque, after compensating the torque losses (i.e. due to coulomb and viscous friction in (3) and using the parameters in Table I) [20, Sec. 11.1.5].

### B. Transient and steady-state responses

The electromechanical transients are directly affected by the *dq* frame dynamics. In order to evaluate the transient response for the proposed *dq* controllers in Sect. III, the reference *dq* currents (i.e.  $i_{s,ref}^{dq}$ ) are varied within a 0.5 s window, while monitoring the current transients during current references steps at 800 RPM. Fig. 7 illustrates the aforementioned current

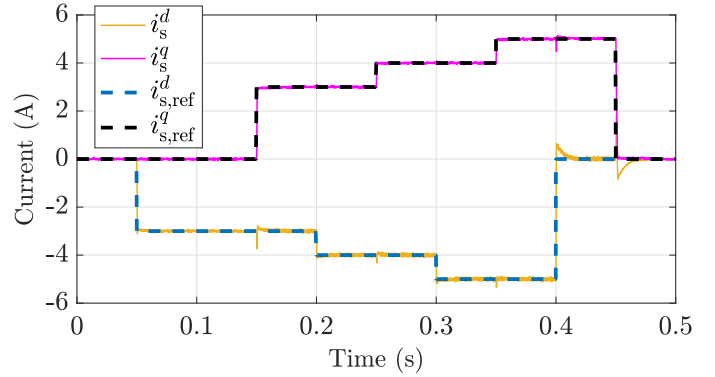


Fig. 7. Experimental transient response for the proposed current controller in the *dq* frame as shown in Sect. III.

references with the corresponding  $i_s^{dq}$ . From Fig. 7, complete decoupling between the *dq* coordinates (i.e. disturbance compensation) could not be achieved; due to the presence of the R-controllers which exhibits some transients during changes in  $i_{s,ref}^{dq}$ . The settling times for a 2% error criteria were 0.62 ms and 1.2 ms for the *dq* coordinates, respectively. The settling times were larger when the reference currents were regulated to zero, which lead to *dq* settling times of 6.8 ms and 8 ms, respectively. The reported results are considered still acceptable.

For the steady-state evaluation, the harmonic spectrum of the stator currents was analyzed through FFT upon applying a given current reference of  $i_{s,ref}^{dq} = (-3.78, -3.78)$  A. The resulted six phase currents are shown in Fig. 8(a). The FFT analysis shown in Fig. 8(b) indicate that the proposed controller is able to suppress the harmonic currents according to the indicated resonant frequencies shown in Fig. 1. A total harmonic distortion (THD) of 0.98% is reported. Furthermore, the measured electromechanical torque is compared with the computed torques leading to the results shown in Fig. 9, indicating the effectiveness of the proposed electromechanical torque model using the non-linear flux linkage maps in (5).

### C. Torque maps and MTPA criteria

The proposed electromechanical model using the flux maps in (5) is evaluated in the entire *dq* space compared to the linear flux linkages assumptions provided by the manufacturer and proposed method in [15, Sec. 5.3]. This is done by forcing  $i_{s,ref}^{dq}$  through the entire *dq* space resulting in a surface plot defined as *torque maps*. By plotting the difference between the measured torque and computed torques, the outcomes are shown in Fig. 10, indicating that the proposed method produces the minimal error compared to the linear flux linkage method. It should be noted that the error between the proposed method and the measured electromechanical torque in Fig. 10(c) is not perfectly zero; since the acquired flux linkage maps in Fig. 2(a)-(b) had undergone some off-line smoothing leading to numerical errors. Furthermore, compensating the effect of torque losses is not perfect over the entire speed range. Nevertheless, the errors in Fig. 10(a)-(b) tend to increase and

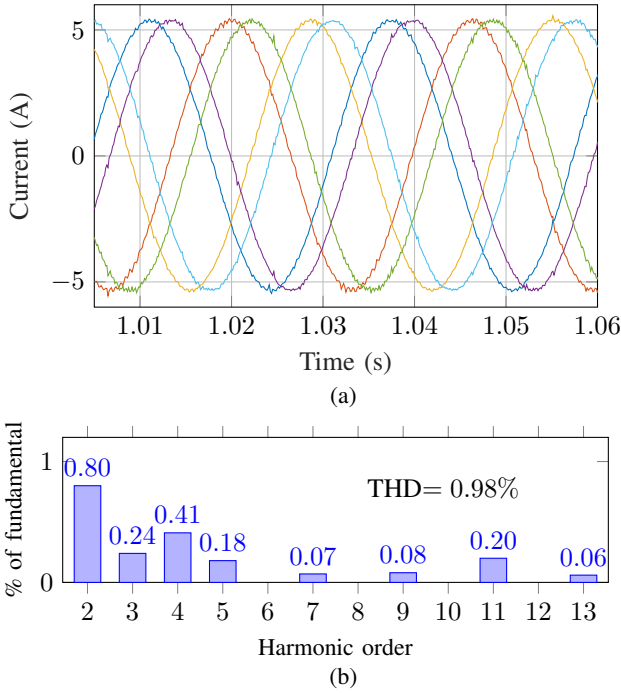


Fig. 8. (a) Stator currents and the corresponding (b) FFT at steady-state for  $\mathbf{i}_{s,ref}^{dq} = (-3.78, -3.78)^\top$  A and  $\mathbf{i}_{s,ref}^{xy} = (0, 0)^\top$  A.

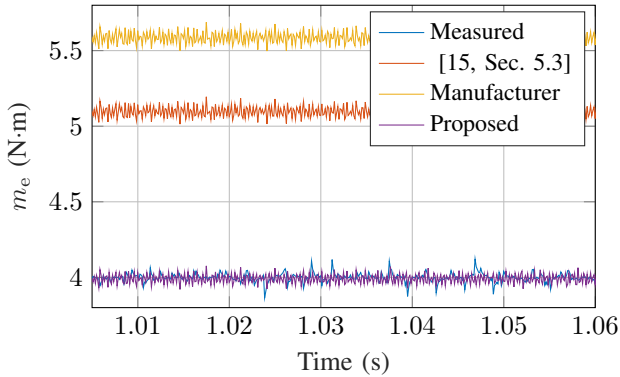


Fig. 9. Comparison between the electromechanical torques computed using the parameters acquired by different methods and the measured one.

vary as the operating point is not in the vicinity of the origin point; since the effect of saturation and coupling between the  $dq$  coordinates starts to prevail.

On the other hand, it is possible to use torque map using the non-linear flux linkage to obtain the MTPA loci, which are obtained by calculating the minimum current magnitude  $\|\mathbf{i}_s^{dq}\|$  for a given measured torque. The new MTPA loci is shown in Fig. 11. It is clear that using non-linear flux linkage maps results in different MTPA loci compared to computing such loci with constant parameters. In other words, the proposed method leads to an optimum MTPA current reference generation compared to the other methods (i.e. reduced copper losses) [10]–[12]. By defining the stator copper losses as  $P_{Cu} := 3\|\mathbf{i}_s^{dq}\|^2 R_s$  (assuming  $\mathbf{i}_s^{xy}$  is regulated to zero), Fig. 12 shows the percentage errors between the stator

copper losses  $\Delta P_{Cu}\%$  produced by the MTPA loci, obtained by the manufacturer’s and [15, Sec. 5.3] parameters, and the non-linear flux linkage maps. Thus, the effectiveness of the proposed method is verified.

## V. CONCLUSION

In this paper, an enhanced electromechanical model was presented for ADT-IPMSMs, relying on estimating the non-linear flux linkages, from which the differential inductances are identified, taking into account the cross-coupling between the  $dq$  coordinates and saturation. The  $dq$  harmonics, due to the PM flux harmonics, are also identified. The PI and R-controllers were tuned and operating-point adaptive depending on the obtained non-linear flux linkages. The non-linear flux linkages affects both the generated electromechanical torque and the MTPA current loci, which have been evaluated against assuming linear flux linkages as given by the manufacturer and estimated by recent literature. Thus, the motivation of employing such non-linear flux linkage maps, in terms of modeling and control has been verified. The theoretical findings were corroborated with a 2.5 kW ADT-IPMSM setup for the sake of validation.

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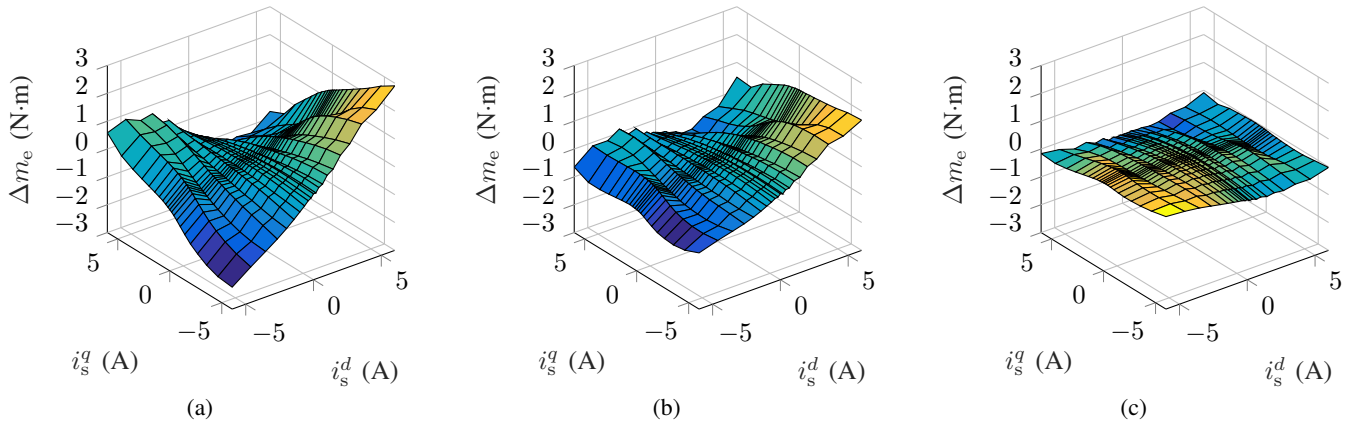


Fig. 10. Error between the measured electromechanical torque and the computed one using the parameters provided by (a) the manufacturer, (b) [15, Sec. 5.3], and (c) the non-linear flux linkage maps.

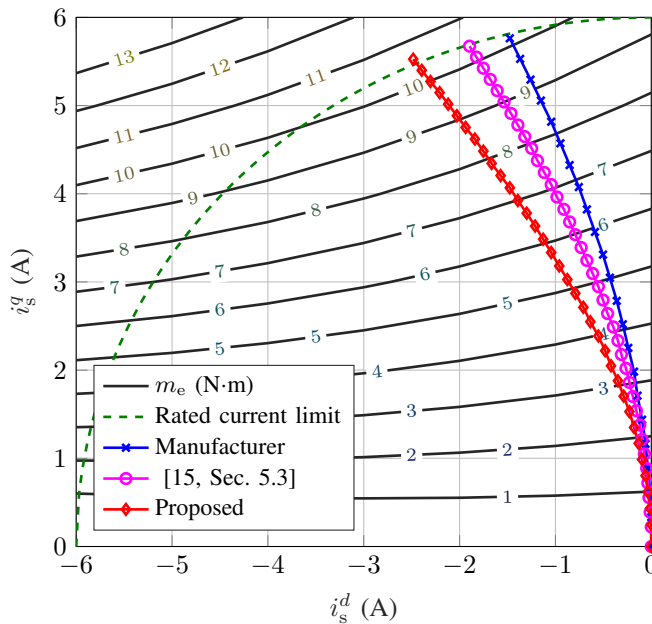


Fig. 11. MTPA current loci for the different torque computation methods and their intersection with the measured torque parabolas.

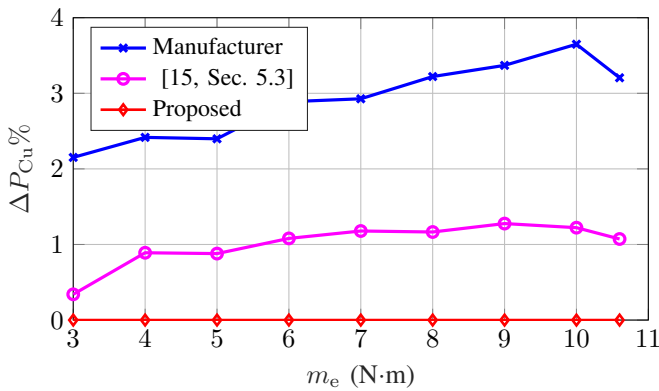


Fig. 12. Comparison of  $\Delta P_{Cu}\%$  computed with the parameters of the different methods with respect to the losses computed when adopting the non-linear flux linkage maps.

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